In The Name of God The Most

Compassionate, The Most Merciful



### **General Theory of Electric Machines**



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### **Induction Motors and Generators**



An induction machine can be used as a motor or a generator.

• Electric motors receive electrical energy as input and provide mechanical energy as output.



• Electric generators receive mechanical energy as input and provide electrical energy as output.







#### 1. Stator

- The stator core is made from the **laminated** steel.
- The armature windings are located in stator slots.
- Three-phase windings are connected in **star** or **delta** configuration.







#### 2. Rotor

#### A. Squirrel Cage

- Aluminium or copper bars are located in the rotor **slots**.
- The bars are short-circuited from both sides using end rings.





#### 2. Rotor

#### **B. Wound Rotor**

- Aluminium or copper windings are located in the rotor **slots**.
- There are three slip-rings and brushes used for energy transfer.





 Synchronous speed. Assume the frequency of the applied voltage to the stator winding is f and the machine has P poles; the synchronous speed is defined as:

$$N_s = \frac{120f}{P}$$
 rpm

$$n_s = \frac{2f}{P}$$
 rps

$$\omega_s = \frac{4\pi f}{P}$$
 rad/s



Synchronous speed is a mechanical quantity.



2. Rotor speed is the speed of the rotor.

#### In motoring mode



Rotor speed is a mechanical quantity.



**3. Slip speed** is the difference between the synchronous speed and the rotor speed:



#### Slip speed is a mechanical quantity.



4. Slip is the slip speed divided by the synchronous speed:



#### Slip is a **dimensionless** quantity.

### How Does an Induction Motor Work?



- Connecting the 3-phase stator windings to a 3-phase AC source flows the current in the stator windings.
- 2. The stator current causes a **rotating magnetic field** with synchronous speed.
- 3. The rotating magnetic field, **induces a voltage** in the rotor bars.
- Since the rotor bars are short-circuited by end-rings, a current flows in the rotor bars.
- 5. The rotor bar current produces **another rotating magnetic field** which rotates with synchronous speed in the same direction as the stator magnetic field.
- 6. Electromagnetic **torque** is developed due to the interaction between two magnetic fields.  $T_{em} = kB_r \times B_s$
- 7. The developed torque can **rotate** the rotor.

### What is Meant by Asynchronous?

- The induced voltage in rotor bars is due to the stator rotating magnetic field.
- Therefore if the rotor rotates with synchronous speed, no voltage is induced in the rotor bars and no torque can be developed.
- Hence, to develop torque, the rotor speed should be different from the synchronous speed.
- It is because induction motors are often called **asynchronous** motors.







### **Operating Modes of Induction Machines**



- The stator windings are connected to a 3-phase ac source.
- The mechanical energy is delivered on the motor shaft.

$$\rightarrow N_s$$
  
 $\rightarrow N_r$ 

$$0 \le N_r \le N_s$$

$$1 \ge S \ge 0$$





### **Operating Modes of Induction Machines**



- The stator windings are connected to a 3-phase ac source.
- The machine is in motoring mode.
- If by using a mechanical mover the rotor speed is increased to above synchronous speed, the machine will be a generator.







#### **Operating Modes of Induction Machines**



- If during the motoring mode the sequence of the applied voltage is changed,
- then the rotating magnetic field will change the direction,
- due to rotor inertia the magnetic field speed is in opposite of the rotor speed.
- The rotor will change the direction of rotation if it is not disconnected from the source.





### **Rotating Magnetic Field**



Consider the following **3-phase AC** machine with **2 poles** and **concentrated winding**. The 3-phase AC currents can be expressed as:



#### **Rotating Magnetic Field**



The rotating magnetomotive force (MMF) can be expressed as

$$\vec{F}(\theta,t) = \vec{F}_a(\theta,t) + \vec{F}_b(\theta,t) + \vec{F}_c(\theta,t)$$

 $\vec{F}(\theta,t) = Ni_a(t)\cos\theta + Ni_b(t)\cos(\theta - 120^\circ) + Ni_c(t)\cos(\theta + 120^\circ)$ 



#### **Rotating Magnetic Field**



The rotating magnetomotive force (MMF) can be expressed as

 $\vec{F}(\theta,t) = Ni_a(t)\cos\theta + Ni_b(t)\cos(\theta - 120^\circ) + Ni_c(t)\cos(\theta + 120^\circ)$ 

$$\vec{F}(\theta, t) = NI_m \cos \omega t \cos \theta$$
$$+ NI_m \cos(\omega t - 120^\circ) \cos(\theta - 120^\circ)$$
$$+ NI_m \cos(\omega t + 120^\circ) \cos(\theta + 120^\circ)$$

 $\vec{F}(\theta,t) = \frac{3}{2}NI_m \cos(\omega t - \theta)$ 

$$\cos a \cos b = \frac{1}{2}\cos(a-b) + \frac{1}{2}\cos(a+b)$$

#### **Electrical & Mechanical Angles**



Consider an AC machine with P poles.



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# **Coil Pitch**



• **Coil pitch**: is the angle between two sides of one armature coil in electrical angle. If the coil pitch is 180 electrical degrees, the coil is a **full-pitch coil**; otherwise it is called **short**- or **chorded**-**pitch coil**.



# **Pitch Factor**

Assume the axial length of the stator is l and the coil pitch is  $\rho$ 



# **Pitch Factor**

• Assume the axial length of the stator is l and the coil pitch is  $\rho$  in electrical angle as shown in the figure  $B_{l}$ 





# **Pitch Factor**

• Since the induced voltage may not be ideal sinusoidal, the pitch factor is expressed for all harmonics

$$k_{pn} = \sin\left(\frac{n\rho}{2}\right)$$

- Why is short-pitch coil used?
- Although short-pitch coil decreases the induced voltage (by about 3%), it decreases the disturbing harmonics significantly (by about 70 to 80%).
- Disturbing harmonics in electric machines are 5 and 7;



\**b'**1

Cz

28

 $\otimes$ 

 $a_1$ 

**b**'<sub>3</sub>

 $20^{\circ}$  20

# **Distributed Windings**

**c'**2

 $(\times$ 

 $b_2$ 

*c*′<sub>3</sub>

**b**<sub>1</sub>

- To make the MMF more sinusoidal, distributed windings are used.
- Consider the following 2-pole single-layer distributed winding AC machine:
- The induced voltage in phase *a* is:

$$\vec{E}_{aa'} = \vec{E}_{aa'1} + \vec{E}_{aa'2} + \vec{E}_{aa'3}$$

$$\vec{E}_{aa'} = E_{rms} \angle -20 + E_{rms} \angle 0 + E_{rms} \angle 20$$

$$\vec{E}_{aa'} = 2.88 E_{rms} \angle 0$$

# **Distribution Factor**

• Distribution factor is defined as:

Induced voltage for distributed winding

Induced voltage for concentrated winding

• In previous example

$$k_d = \frac{2.88E_{rms}}{3E_{rms}}$$

$$k_d = 0.96$$



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 $k_d =$ 

# **Multi-layer Winding**



- Most of AC machines, especially large machines, have 2-layer winding.
- 2-layer winding has the following advantages compared to single-layer winding:
  - 1. Simpler winding
  - 2. Simpler connection
  - 3. Higher mechanical strength
  - 4. Optimal use of slots
  - 5. Lower cost





# **Distribution Factor**

• Distribution factor is defined as:

Induced voltage for distributed winding

Induced voltage for concentrated winding

 Assume the angle between two adjacent slots is γ in electrical angle and the number of slots per pole per phase is m, therefore the distribution factor is written as

$$k_d = \frac{\sin\frac{m\gamma}{2}}{m\sin\frac{\gamma}{2}}$$

• In *n*-th harmonic it is

$$k_{dn} = \frac{\sin \frac{mn \gamma}{2}}{m \sin \frac{n \gamma}{2}}$$



 $k_d =$ 

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# Winding Factor

• Winding factor is defined as

 $k_{\rm w} = k_{\rm w} k_{\rm d}$ 

$$k_{p} = \sin\left(\frac{\rho}{2}\right)$$

$$k_{d} = \frac{\sin\frac{mn\gamma}{2}}{m\sin\frac{\gamma}{2}}$$

$$k_{dn} = \frac{\sin\frac{mn\gamma}{2}}{m\sin\frac{n\gamma}{2}}$$

$$k_{dn} = \frac{\sin\frac{mn\gamma}{2}}{m\sin\frac{n\gamma}{2}}$$

$$k_{dn} = \frac{\sin\frac{mn\gamma}{2}}{m\sin\frac{n\gamma}{2}}$$



• **Stator voltage** equations in the **stator** fixed frame:

 $v_{as} = r_s i_{as} + d\lambda_{as} / dt$  $v_{bs} = r_s i_{bs} + d\lambda_{bs} / dt$  $v_{cs} = r_s i_{cs} + d\lambda_{cs} / dt$ 

Rotor voltage equations in the rotor fixed frame:

$$v_{ar} = r_r i_{ar} + d\lambda_{ar}/dt$$
$$v_{br} = r_r i_{br} + d\lambda_{br}/dt$$
$$v_{cr} = r_r i_{cr} + d\lambda_{cr}/dt$$



where  $r_s$  is the per-phase stator resistance and  $r_r$  is the per-phase rotor resistance.



#### 2. Flux linkage equations

• In matrix form, the flux linkages of the stator and rotor windings in terms of the winding inductances and currents are expressed as

$$\begin{bmatrix} \boldsymbol{\lambda}_{s}^{abc} \\ \boldsymbol{\lambda}_{r}^{abc} \end{bmatrix} = \begin{bmatrix} \mathbf{L}_{ss}^{abc} & \mathbf{L}_{sr}^{abc} \\ \mathbf{L}_{rs}^{abc} & \mathbf{L}_{rr}^{abc} \end{bmatrix} \begin{bmatrix} \mathbf{i}_{s}^{abc} \\ \mathbf{i}_{r}^{abc} \end{bmatrix}$$

where





#### 2. Flux linkage equations

• The stator-to-stator inductance matrix is expressed as:

$$\mathbf{L}_{ss}^{abc} = \begin{bmatrix} L_{ls} + L_{ss} & L_{sm} & L_{sm} \\ L_{sm} & L_{ls} + L_{ss} & L_{sm} \\ L_{sm} & L_{sm} & L_{ls} + L_{ss} \end{bmatrix}$$



#### where

- $L_{ls}$  is the per-phase stator winding leakage inductance
- $L_{ss}$  is the per-phase stator winding self-inductance
- $L_{sm}$  is the **mutual-inductance** between **stator** windings.



• The **rotor-to-rotor inductance** matrix is expressed as:

$$\mathbf{L}_{rr}^{abc} = \begin{bmatrix} L_{lr} + L_{rr} & L_{rm} & L_{rm} \\ L_{rm} & L_{lr} + L_{rr} & L_{rm} \\ L_{rm} & L_{rm} & L_{lr} + L_{rr} \end{bmatrix}$$

where

 $L_{lr}$  is the per-phase **rotor** winding **leakage** inductance

- $L_{rr}$  is the per-phase **rotor** winding **self-inductance**
- $L_{rm}$  is the **mutual-inductance** between **rotor** windings.


#### 2. Flux linkage equations

• The stator-to-rotor mutual inductance matrix (which depends on the rotor angle) is expressed as:

$$\mathbf{L}_{sr}^{abc} = \begin{bmatrix} \mathbf{L}_{rs}^{abc} \end{bmatrix}^{T} = L_{sr} \begin{bmatrix} \cos\theta_{r} & \cos(\theta_{r} + 2\pi/3) & \cos(\theta_{r} - 2\pi/3) \\ \cos(\theta_{r} - 2\pi/3) & \cos\theta_{r} & \cos(\theta_{r} + 2\pi/3) \\ \cos(\theta_{r} + 2\pi/3) & \cos(\theta_{r} - 2\pi/3) & \cos\theta_{r} \end{bmatrix}$$

#### where

*L<sub>sr</sub>* is the **peak** value of the **stator-to-rotor mutual** inductance

 $\theta_r$  is the rotor angle w.r.t. the stator fixed frame

$$\theta_r = \omega_r t + \theta_{r0}$$



#### 2. Flux linkage equations

• If the reluctive drops in iron are neglected, the inductances can be expressed in terms of the air-gap permeance  $(P_g)$ , the rotor winding turns  $(N_r)$  and the stator winding turns  $(N_s)$ :

$$L_{ss} = N_s^2 P_g \qquad L_{rr} = N_r^2 P_g$$
$$L_{sm} = N_s^2 P_g \cos(2\pi/3)$$
$$L_{rm} = N_r^2 P_g \cos(2\pi/3)$$

 $L_{sr} = N_s N_r P_g$ 





$$\begin{bmatrix} v_{as} \\ v_{bs} \\ v_{cs} \\ v_{ar} \\ v_{br} \\ v_{cr} \end{bmatrix} = \begin{bmatrix} r_s + pL_s & pL_{sm} & pL_{sm} & pL_{sr}\cos\theta_r & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_2 \\ pL_{sm} & r_s + pL_s & pL_{sm} & pL_{sr}\cos\theta_2 & pL_{sr}\cos\theta_r & pL_{sr}\cos\theta_1 \\ pL_{sm} & pL_{sm} & r_s + pL_s & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_2 & pL_{sr}\cos\theta_r \\ pL_{sr}\cos\theta_r & pL_{sr}\cos\theta_2 & pL_{sr}\cos\theta_1 & r_r + pL_r & pL_{rm} & pL_{rm} \\ pL_{sr}\cos\theta_r & pL_{sr}\cos\theta_r & pL_{sr}\cos\theta_2 & pL_{rm} & r_r + pL_r & pL_{rm} \\ pL_{sr}\cos\theta_2 & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r & pL_{sr}\cos\theta_r & pL_{rm} \\ pL_{sr}\cos\theta_2 & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r & pL_{rm} & r_r + pL_r \\ pL_{sr}\cos\theta_2 & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r & pL_{rm} & pL_{rm} \\ pL_{sr}\cos\theta_2 & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r & pL_{rm} & pL_{rm} \\ pL_{sr}\cos\theta_2 & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r & pL_{rm} & pL_{rm} \\ pL_{sr}\cos\theta_2 & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r & pL_{rm} & pL_{rm} \\ pL_{sr}\cos\theta_2 & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r & pL_{rm} & pL_{rm} \\ pL_{sr}\cos\theta_2 & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r & pL_{rm} \\ pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r \\ pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r & pL_{sr}\cos\theta_r \\ pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r \\ pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r \\ pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r \\ pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r & pL_{sr}\cos\theta_r \\ pL_{sr}\cos\theta_1 & pL_{sr}\cos\theta_r \\ pL$$

where 
$$\theta_1 = \theta_r + 2\pi/3$$
  $\theta_2 = \theta_r - 2\pi/3$   $p = d/dt$   
 $L_s = L_{ls} + L_{ss}$   $L_r = L_{lr} + L_{rr}$ 

• It can be written in the compact matrix form as follows

$$\begin{bmatrix} \mathbf{v}_{s}^{abc} \\ \mathbf{v}_{r}^{abc} \end{bmatrix} = \begin{bmatrix} \mathbf{r}_{s}^{abc} & \mathbf{0} \\ \mathbf{0} & \mathbf{r}_{r}^{abc} \end{bmatrix} \begin{bmatrix} \mathbf{i}_{s}^{abc} \\ \mathbf{i}_{r}^{abc} \end{bmatrix} + \frac{d}{dt} \begin{bmatrix} \mathbf{L}_{ss}^{abc} & \mathbf{L}_{sr}^{abc} \\ \mathbf{L}_{rs}^{abc} & \mathbf{L}_{rr}^{abc} \end{bmatrix} \begin{bmatrix} \mathbf{i}_{s}^{abc} \\ \mathbf{i}_{r}^{abc} \end{bmatrix}$$

where







- As mentioned before the machine model in *abc* system is represented by **six first-order** differential equations.
- These differential equations are **coupled** due to the mutual inductances.
- The stator-to-rotor coupling terms are functions of rotor position; thus when rotor rotates, they vary with time.
- Mathematical transformations like dq or  $\alpha\beta$  can facilitate the computation by transferring the differential equations with time-varying inductances to DEs with constant inductances.



- The equations of IM are first derived in the arbitrary reference frame (RF) rotating at a speed of ω in the direction of the rotor rotation.
- Setting ω = 0 yields the equations in the stationary RF.
- Setting  $\omega = \omega_e$  yields the equations in the synchronous RF.





#### The transformation is expressed as **Arbitrary** q ω reference frame $\left[\mathbf{f}_{qd0}\right] = \left[\mathbf{T}_{qd0}(\boldsymbol{\theta})\right] \left[\mathbf{f}_{abc}\right]$ $\omega_r$ ) as θ • The transformation angle is Rotor -CS -bs ) ar fixed $\bigotimes$ ∕⊗ -br $\theta_r$ frame $\theta(t) = \int_0^t \omega(t) \, dt + \theta(0)$ ⊗-cr Stator cr fixed bs -ar br frame $\bigcirc$ The rotor angle is expressed as d CS -as $\otimes$ $\theta_r(t) = \int_0^t \omega_r(t) dt + \theta_r(0)$

• The transformation and its inverse for **stator** equations

$$\begin{bmatrix} \mathbf{f}_{qd0} \end{bmatrix} = \begin{bmatrix} \mathbf{T}_{qd0}(\theta) \end{bmatrix} \begin{bmatrix} \mathbf{f}_{abc} \end{bmatrix} \qquad \begin{bmatrix} \mathbf{f}_{abc} \end{bmatrix} = \begin{bmatrix} \mathbf{T}_{qd0}(\theta) \end{bmatrix}^{-1} \begin{bmatrix} \mathbf{f}_{qd0} \end{bmatrix}$$
$$\begin{bmatrix} \mathbf{T}_{qd0}(\theta) \end{bmatrix} = \frac{2}{3} \begin{bmatrix} \cos\theta & \cos(\theta - 2\pi/3) & \cos(\theta + 2\pi/3) \\ \sin\theta & \sin(\theta - 2\pi/3) & \sin(\theta + 2\pi/3) \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix}$$
$$\begin{bmatrix} \mathbf{T}_{qd0}(\theta) \end{bmatrix}^{-1} = \begin{bmatrix} \cos\theta & \sin\theta & 1 \\ \cos(\theta - 2\pi/3) & \sin(\theta - 2\pi/3) & 1 \\ \cos(\theta + 2\pi/3) & \sin(\theta + 2\pi/3) & 1 \end{bmatrix}$$



$$\begin{bmatrix} \mathbf{f}_{qd0} \end{bmatrix} = \begin{bmatrix} \mathbf{T}_{qd0}(\theta - \theta_r) \end{bmatrix} \begin{bmatrix} \mathbf{f}_{abc} \end{bmatrix} \qquad \begin{bmatrix} \mathbf{f}_{abc} \end{bmatrix} = \begin{bmatrix} \mathbf{T}_{qd0}(\theta - \theta_r) \end{bmatrix}^{-1} \begin{bmatrix} \mathbf{f}_{qd0} \end{bmatrix}$$
$$\begin{bmatrix} \mathbf{T}_{qd0}(\theta - \theta_r) \end{bmatrix} = \frac{2}{3} \begin{bmatrix} \cos(\theta - \theta_r) & \cos(\theta - \theta_r - 2\pi/3) & \cos(\theta - \theta_r + 2\pi/3) \\ \sin(\theta - \theta_r) & \sin(\theta - \theta_r - 2\pi/3) & \sin(\theta - \theta_r + 2\pi/3) \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix}$$

$$\begin{bmatrix} \mathbf{T}_{qd0}(\theta - \theta_r) \end{bmatrix}^{-1} = \begin{bmatrix} \cos(\theta - \theta_r) & \sin(\theta - \theta_r) & 1 \\ \cos(\theta - \theta_r - 2\pi/3) & \sin(\theta - \theta_r - 2\pi/3) & 1 \\ \cos(\theta - \theta_r + 2\pi/3) & \sin(\theta - \theta_r + 2\pi/3) & 1 \end{bmatrix}$$

qd0 voltage equations

Stator

• The stator voltage equations in *abc* system is expressed as:

 $\mathbf{v}_s^{abc} = \mathbf{r}_s^{abc} \mathbf{i}_s^{abc} + p \boldsymbol{\lambda}_s^{abc}$ 

where 
$$p = d/dt$$

• Applying the transformation  $\mathbf{T}_{qd0}(\theta)$  to voltage, current and flux yields  $[\mathbf{f}_{abc}] = [\mathbf{T}_{qd0}(\theta)]^{-1} [\mathbf{f}_{qd0}]$ 

$$\left[\mathbf{T}_{qd0}(\boldsymbol{\theta})\right]^{-1}\mathbf{v}_{s}^{qd0} = \mathbf{r}_{s}^{abc}\left[\mathbf{T}_{qd0}(\boldsymbol{\theta})\right]^{-1}\mathbf{i}_{s}^{qd0} + p\left(\left[\mathbf{T}_{qd0}(\boldsymbol{\theta})\right]^{-1}\boldsymbol{\lambda}_{s}^{qd0}\right)\right)$$

• Multiplying both sides by the transformation matrix yields

$$\mathbf{v}_{s}^{qd0} = \left[\mathbf{T}_{qd0}(\theta)\right]\mathbf{r}_{s}^{abc}\left[\mathbf{T}_{qd0}(\theta)\right]^{-1}\mathbf{i}_{s}^{qd0} + \left[\mathbf{T}_{qd0}(\theta)\right]p\left(\left[\mathbf{T}_{qd0}(\theta)\right]^{-1}\boldsymbol{\lambda}_{s}^{qd0}\right)\right)$$

qd0 voltage equations

 $\mathbf{v}_{s}^{qd0} = \left[\mathbf{T}_{qd0}(\theta)\right]\mathbf{r}_{s}^{abc}\left[\mathbf{T}_{qd0}(\theta)\right]^{-1}\mathbf{i}_{s}^{qd0} + \left[\mathbf{T}_{qd0}(\theta)\right]p\left(\left[\mathbf{T}_{qd0}(\theta)\right]^{-1}\boldsymbol{\lambda}_{s}^{qd0}\right)\right)$ 

Substituting the following relations in the above expression

$$\mathbf{r}_{s}^{abc} = \mathbf{r}_{s}^{qd0} = r_{s} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \mathbf{T}_{qd0}(\theta) \end{bmatrix} p \left( \begin{bmatrix} \mathbf{T}_{qd0}(\theta) \end{bmatrix}^{-1} \lambda_{s}^{qd0} \right) = \omega \begin{bmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \lambda_{s}^{qd0} + p \lambda_{s}^{qd0}$$

yields

$$\mathbf{v}_{s}^{qd0} = \mathbf{r}_{s}^{qd0} \mathbf{i}_{s}^{qd0} + \omega \begin{bmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \boldsymbol{\lambda}_{s}^{qd0} + p \,\boldsymbol{\lambda}_{s}^{qd0}$$

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Stator





qd0 voltage equations

Rotor

• The **rotor voltage** equations in *abc* system is expressed as:

 $\mathbf{v}_r^{abc} = \mathbf{r}_r^{abc} \mathbf{i}_r^{abc} + p \boldsymbol{\lambda}_r^{abc}$ 

• Similarly applying the transformation  $\mathbf{T}_{qd0}(\theta - \theta_r)$  to voltage, current and flux yields  $[\mathbf{f}_{abc}] = [\mathbf{T}_{qd0}(\theta - \theta_r)]^{-1} [\mathbf{f}_{qd0}]$ 

$$\mathbf{v}_{r}^{qd0} = \mathbf{r}_{r}^{qd0} \mathbf{i}_{r}^{qd0} + (\omega - \omega_{r}) \begin{bmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \lambda_{r}^{qd0} + p \lambda_{r}^{qd0}$$



![](_page_51_Figure_0.jpeg)

#### **Induction Machine Modelling in** Arbitrary *qd*0 RF **Stator & Rotor**

![](_page_52_Figure_1.jpeg)

![](_page_52_Figure_2.jpeg)

![](_page_53_Picture_1.jpeg)

![](_page_53_Picture_2.jpeg)

• The stator flux linkage equations in *abc* system is expressed as:

 $\boldsymbol{\lambda}_{s}^{abc} = \mathbf{L}_{ss}^{abc} \, \mathbf{i}_{s}^{abc} + \mathbf{L}_{sr}^{abc} \, \mathbf{i}_{r}^{abc}$ 

• Applying the transformation  $\mathbf{T}_{qd0}(\theta)$  and  $\mathbf{T}_{qd0}(\theta - \theta_r)$  to the stator and rotor quantities yields

$$\left[\mathbf{T}_{qd0}(\theta)\right]^{-1}\boldsymbol{\lambda}_{s}^{qd0} = \mathbf{L}_{ss}^{abc}\left[\mathbf{T}_{qd0}(\theta)\right]^{-1}\mathbf{i}_{s}^{qd0} + \mathbf{L}_{sr}^{abc}\left[\mathbf{T}_{qd0}(\theta - \theta_{r})\right]^{-1}\mathbf{i}_{r}^{qd0}$$

• Multiplying both sides by  $\mathbf{T}_{qd0}(\theta)$  yields

$$\boldsymbol{\lambda}_{s}^{qd0} = \left[\mathbf{T}_{qd0}(\boldsymbol{\theta})\right] \mathbf{L}_{ss}^{abc} \left[\mathbf{T}_{qd0}(\boldsymbol{\theta})\right]^{-1} \mathbf{i}_{s}^{qd0} + \left[\mathbf{T}_{qd0}(\boldsymbol{\theta})\right] \mathbf{L}_{sr}^{abc} \left[\mathbf{T}_{qd0}(\boldsymbol{\theta} - \boldsymbol{\theta}_{r})\right]^{-1} \mathbf{i}_{r}^{qd0}$$

qd0 flux linkage relation

$$\boldsymbol{\lambda}_{s}^{qd0} = \begin{bmatrix} \mathbf{T}_{qd0}(\theta) \end{bmatrix} \mathbf{L}_{ss}^{abc} \begin{bmatrix} \mathbf{T}_{qd0}(\theta) \end{bmatrix}^{-1} \mathbf{i}_{s}^{qd0} + \begin{bmatrix} \mathbf{T}_{qd0}(\theta) \end{bmatrix} \mathbf{L}_{sr}^{abc} \begin{bmatrix} \mathbf{T}_{qd0}(\theta - \theta_{r}) \end{bmatrix}^{-1} \mathbf{i}_{r}^{qd0}$$
$$\mathbf{L}_{ss}^{qd0} = \mathbf{L}_{ss}^{qd0} \mathbf{L}_{sr}^{qd0}$$

$$\lambda_s^{qd0} = \mathbf{L}_{ss}^{qd0} \, \mathbf{i}_s^{qd0} + \mathbf{L}_{sr}^{qd0} \, \mathbf{i}_r^{qd0}$$

The inductance matrices  $L_{ss}$  and  $L_{sr}$  in *qd*0 reference frame need to be obtained.

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**Stator** 

qd0 flux linkage relation

Stator

where  $\mathbf{L}_{ss}^{qd0} = [\mathbf{T}_{qd0}(\theta)]\mathbf{L}_{ss}^{abc}[\mathbf{T}_{qd0}(\theta)]^{-1}$  is the stator self-inductance

matrix in qd0 reference frame

$$\mathbf{L}_{ss}^{qd0} = \frac{2}{3} \begin{bmatrix} \cos\theta & \cos\left(\theta - \frac{2\pi}{3}\right) & \cos\left(\theta + \frac{2\pi}{3}\right) \\ \sin\theta & \sin\left(\theta - \frac{2\pi}{3}\right) & \sin\left(\theta + \frac{2\pi}{3}\right) \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix} \begin{bmatrix} L_{ls} + L_{ss} & L_{sm} & L_{ls} + L_{ss} \\ L_{sm} & L_{ls} + L_{ss} & L_{sm} \\ L_{sm} & L_{sm} & L_{ls} + L_{ss} \end{bmatrix} \begin{bmatrix} \cos\theta & \sin\theta & 1 \\ \cos\left(\theta - \frac{2\pi}{3}\right) & \sin\left(\theta - \frac{2\pi}{3}\right) & 1 \\ \cos\left(\theta + \frac{2\pi}{3}\right) & \sin\left(\theta + \frac{2\pi}{3}\right) & 1 \end{bmatrix}$$

$$\mathbf{L}_{ss}^{qd0} = \begin{bmatrix} L_{ls} + \frac{3}{2}L_{ss} & 0 & 0\\ 0 & L_{ls} + \frac{3}{2}L_{ss} & 0\\ 0 & 0 & L_{ls} \end{bmatrix}$$
Diagonal matrix and independent of  $\theta$ .

qd0 flux linkage relation

Stator

 $(-) \neg$ 

and  $\mathbf{L}_{sr}^{qd0} = [\mathbf{T}_{qd0}(\theta)] \mathbf{L}_{sr}^{abc} [\mathbf{T}_{qd0}(\theta - \theta_r)]^{-1}$ 

is the stator-to-rotor

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mutual inductance matrix.

where

$$\mathbf{L}_{sr}^{abc} = L_{sr} \begin{bmatrix} \cos\theta_r & \cos(\theta_r + 2\pi/3) & \cos(\theta_r - 2\pi/3) \\ \cos(\theta_r - 2\pi/3) & \cos\theta_r & \cos(\theta_r + 2\pi/3) \\ \cos(\theta_r + 2\pi/3) & \cos(\theta_r - 2\pi/3) & \cos\theta_r \end{bmatrix}$$
$$\mathbf{L}_{sr}^{qd0} = \begin{bmatrix} \frac{3}{2}L_{sr} & 0 & 0 \\ 0 & \frac{3}{2}L_{sr} & 0 \\ 0 & 0 & 0 \end{bmatrix}$$
$$\mathbf{D}_{iagonal matrix and independent of } \theta.$$

10

(-)

![](_page_57_Picture_1.jpeg)

qd0 flux linkage relation

![](_page_57_Picture_3.jpeg)

• The rotor flux linkage equations in *abc* system is expressed as:

 $\boldsymbol{\lambda}_{r}^{abc} = \mathbf{L}_{rs}^{abc} \, \mathbf{i}_{s}^{abc} + \mathbf{L}_{rr}^{abc} \, \mathbf{i}_{r}^{abc}$ 

• Applying the transformation  $\mathbf{T}_{qd0}(\theta)$  and  $\mathbf{T}_{qd0}(\theta - \theta_r)$  to the stator and rotor quantities yields

$$\left[\mathbf{T}_{qd0}(\theta - \theta_r)\right]^{-1} \boldsymbol{\lambda}_r^{qd0} = \mathbf{L}_{rs}^{abc} \left[\mathbf{T}_{qd0}(\theta)\right]^{-1} \mathbf{i}_s^{qd0} + \mathbf{L}_{rr}^{abc} \left[\mathbf{T}_{qd0}(\theta - \theta_r)\right]^{-1} \mathbf{i}_r^{qd0}$$

• Multiplying both sides by  $\mathbf{T}_{qd0}(\theta - \theta_r)$  yields

$$\boldsymbol{\lambda}_{r}^{qd0} = \left[\mathbf{T}_{qd0}(\boldsymbol{\theta} - \boldsymbol{\theta}_{r})\right] \mathbf{L}_{rs}^{abc} \left[\mathbf{T}_{qd0}(\boldsymbol{\theta})\right]^{-1} \mathbf{i}_{s}^{qd0} + \left[\mathbf{T}_{qd0}(\boldsymbol{\theta} - \boldsymbol{\theta}_{r})\right] \mathbf{L}_{rr}^{abc} \left[\mathbf{T}_{qd0}(\boldsymbol{\theta} - \boldsymbol{\theta}_{r})\right]^{-1} \mathbf{i}_{r}^{qd0}$$

qd0 flux linkage relation

![](_page_58_Figure_2.jpeg)

$$\lambda_r^{qd0} = \mathbf{L}_{rs}^{qd0} \,\mathbf{i}_s^{qd0} + \mathbf{L}_{rr}^{qd0} \,\mathbf{i}_r^{qd0}$$

The inductance matrices  $L_{rs}$  and  $L_{rr}$  in *qd*0 reference frame need to be obtained.

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Rotor

qd0 flux linkage relation

![](_page_59_Picture_2.jpeg)

where  $\mathbf{L}_{rr}^{qd0} = \left[\mathbf{T}_{qd0}(\theta - \theta_r)\right] \mathbf{L}_{rr}^{abc} \left[\mathbf{T}_{qd0}(\theta - \theta_r)\right]^{-1}$  is the rotor self-

inductance matrix in *qd*0 reference frame

$$\mathbf{L}_{rr}^{qd0} = \begin{bmatrix} L_{lr} + \frac{3}{2}L_{rr} & 0 & 0\\ 0 & L_{lr} + \frac{3}{2}L_{rr} & 0\\ 0 & 0 & L_{lr} \end{bmatrix}$$

Diagonal matrix and independent of  $\theta$ .

qd0 flux linkage relation

![](_page_60_Picture_2.jpeg)

and 
$$\mathbf{L}_{rs}^{qd0} = \left[\mathbf{T}_{qd0}(\theta - \theta_r)\right] \mathbf{L}_{rs}^{abc} \left[\mathbf{T}_{qd0}(\theta)\right]^{-1}$$

is the rotor-to-stator

mutual inductance matrix.

$$\mathbf{L}_{rs}^{qd0} = \begin{bmatrix} \frac{3}{2}L_{sr} & 0 & 0\\ 0 & \frac{3}{2}L_{sr} & 0\\ 0 & 0 & 0 \end{bmatrix}$$

Diagonal matrix and independent of  $\theta$ .

*qd0* flux linkage relation

**Stator & Rotor** 

Therefore the flux linkage relation of induction machines in qd0 reference frame is expressed as

$\left\lceil \lambda_{qs}  ight ceil$		$\int L_{ls} + \frac{3}{2} L_{ss}$	0	0	$\frac{3}{2}L_{sr}$	0	0	$\begin{bmatrix} i_{qs} \end{bmatrix}$
$\lambda_{ds}$		0	$L_{ls} + \frac{3}{2}L_{ss}$	0	0	$\frac{3}{2}L_{sr}$	0	$i_{ds}$
$\lambda_{0s}$		0	0	$L_{ls}$	0	0	0	$i_{0s}$
$\lambda_{qr}$	_	$\frac{3}{2}L_{sr}$	0	0	$L_{lr} + \frac{3}{2}L_{rr}$	0	0	$\overline{i_{qr}}$
$\lambda_{dr}$		0	$\frac{3}{2}L_{sr}$	0	0	$L_{lr} + \frac{3}{2}L_{rr}$	0	<i>i</i> <sub>dr</sub>
$\lfloor \lambda_{0r} \rfloor$		0	0	0	0	0	$L_{lr}$	$\lfloor i_{0r} \rfloor$

It is required to refer all rotor quantities to the stator side.

![](_page_62_Picture_1.jpeg)

![](_page_62_Figure_3.jpeg)

The magnetizing inductance,  $L_m$ , on the stator side is expressed as

$$L_{m} = \frac{3}{2} L_{ss} = \frac{3}{2} \frac{N_{s}}{N_{r}} L_{sr} = \frac{3}{2} \left(\frac{N_{s}}{N_{r}}\right)^{2} L_{rr}$$

![](_page_62_Picture_6.jpeg)

*qd0* flux linkage relation referred to stator side

**Stator & Rotor** 

Therefore *qd*0 flux linkage relation is expressed as

$$\begin{bmatrix} \lambda_{qs} \\ \lambda_{ds} \\ \lambda_{ds} \\ \frac{\lambda_{0s}}{\lambda_{qr}} \\ \lambda'_{qr} \\ \lambda'_{dr} \\ \lambda'_{0r} \end{bmatrix} = \begin{bmatrix} L_{ls} + L_m & 0 & 0 & L_m & 0 & 0 \\ 0 & L_{ls} + L_m & 0 & 0 & L_m & 0 \\ 0 & 0 & L_{ls} & 0 & 0 & 0 \\ L_m & 0 & 0 & L'_{lr} + L_m & 0 & 0 \\ 0 & L_m & 0 & 0 & L'_{lr} + L_m & 0 \\ 0 & 0 & 0 & 0 & 0 & L'_{lr} + L_m & 0 \\ 0 & 0 & 0 & 0 & 0 & L'_{lr} \end{bmatrix} \begin{bmatrix} i_{qs} \\ i_{ds} \\ i_{ds} \\ i_{ds} \\ i'_{dr} \\ i'_{dr} \\ i'_{or} \end{bmatrix}$$

where

 $\frac{N_s}{N}$ 

![](_page_63_Picture_11.jpeg)

![](_page_64_Picture_1.jpeg)

• Lets prove the following relation

$$L_m = \frac{3}{2} L_{ss} = \frac{3}{2} \frac{N_s}{N_r} L_{sr} = \frac{3}{2} \left(\frac{N_s}{N_r}\right)^2 L_{rr}$$

• Remind the aforementioned relations

$$L_{ss} = N_s^2 P_g \qquad \qquad L_{rr} = N_r^2 P_g \qquad \qquad L_{sr} = N_s N_r P_g$$

• Therefore using the above relations:

$$L_{ss} = N_s^2 P_g = \left(\frac{N_s}{N_r}\right)^2 N_r^2 P_g = \left(\frac{N_s}{N_r}\right)^2 L_{rr}$$

$$L_{ss} = N_s^2 P_g = \left(\frac{N_s}{N_r}\right) N_s N_r P_g = \left(\frac{N_s}{N_r}\right) L_{sr}$$

![](_page_65_Picture_1.jpeg)

Therefore qd0 model is expressed as

![](_page_65_Figure_3.jpeg)

![](_page_65_Figure_4.jpeg)

• Using the **reactances** instead of the **inductances** in the **stator** voltage and flux linkage equations yields:

$$\begin{aligned}
\psi_{qs} &= r_{s} \, \dot{i}_{qs} + \frac{1}{\omega_{b}} \frac{d\psi_{qs}}{dt} + \frac{\omega}{\omega_{b}} \psi_{ds} \\
\psi_{qs} &= \omega_{b} \, \lambda_{qs} \\
\psi_{ds} &= r_{s} \, \dot{i}_{ds} + \frac{1}{\omega_{b}} \frac{d\psi_{ds}}{dt} - \frac{\omega}{\omega_{b}} \psi_{qs} \\
\psi_{ds} &= \omega_{b} \, \lambda_{ds} \\
\psi_{ds} &= \omega_{b} \, \lambda_{ds} \\
\psi_{ds} &= \omega_{b} \, L_{ls} \\
\psi_{ds} &= \omega_{b} \, L_{ls} \\
\psi_{ds} &= \omega_{b} \, L_{m}
\end{aligned}$$

![](_page_67_Picture_1.jpeg)

• Similarly using the **reactances** instead of the **inductances** in the **rotor** voltage and flux linkage equations (referred to stator) yields:

$$\psi'_{qr} = r'_{r}i'_{qr} + \frac{1}{\omega_{b}}\frac{d\psi'_{qr}}{dt} + \frac{\omega - \omega_{r}}{\omega_{b}}\psi'_{dr}$$

$$\psi'_{qr} = \omega_{b}\lambda'_{qr}$$

$$\psi'_{qr} = \omega_{b}\lambda'_{dr}$$

$$\psi'_{dr} = \omega_{b}\lambda'_{dr}$$

$$\psi'_{dr} = \omega_{b}\lambda'_{dr}$$

$$\psi'_{dr} = \omega_{b}\lambda'_{dr}$$

$$\chi'_{lr} = \omega_{b}L'_{lr}$$

$$\chi'_{m} = \omega_{b}L_{m}$$

q-axis equivalent circuit

$$\psi_{qs} = x_{ls} \, i_{qs} + x_m (i_{qs} + i'_{qr})$$

**Stator** 

$$v'_{qr} = r'_r i'_{qr} + \frac{1}{\omega_b} \frac{d\psi'_{qr}}{dt} + \frac{\omega - \omega_r}{\omega_b} \psi'_{dr}$$

 $v_{qs} = r_s i_{qs} + \frac{1}{\omega_b} \frac{d\psi_{qs}}{dt} + \frac{\omega}{\omega_b} \psi_{ds}$ 

$$\psi'_{qr} = x'_{lr} i'_{qr} + x_m (i_{qs} + i'_{qr})$$

![](_page_68_Figure_6.jpeg)

![](_page_69_Picture_1.jpeg)

Stator 
$$v_{ds} = r_s i_{ds} + \frac{1}{\omega_b} \frac{d\psi_{ds}}{dt} - \frac{\omega}{\omega_b} \psi_{qs}$$
  $\psi_{ds} = x_{ls} i_{ds} + x_m (i_{ds} + i'_{dr})$   
Rotor  $v'_{dr} = r'_r i'_{dr} + \frac{1}{\omega_b} \frac{d\psi'_{dr}}{dt} - \frac{\omega - \omega_r}{\omega_b} \psi'_{qr}$   $\psi'_{dr} = x'_{lr} i'_{dr} + x_m (i_{ds} + i'_{dr})$ 

![](_page_69_Figure_3.jpeg)

0 component equivalent circuit

![](_page_70_Figure_2.jpeg)

#### Induction Machine Modelling in Arbitrary qd0 RF Torque Equation

• The instantaneous input power in *abc* system is obtained as

 $P_{in} = v_{as} \, i_{as} + v_{bs} \, i_{bs} + v_{cs} \, i_{cs} + v'_{ar} \, i'_{ar} + v'_{br} \, i'_{br} + v'_{cr} \, i'_{cr}$ 

• In terms of qd0 quantities, the instantaneous input power is

$$P_{in} = \frac{3}{2} \left( v_{qs} \, i_{qs} + v_{ds} \, i_{ds} + 2v_{0s} \, i_{0s} + v'_{qr} \, i'_{qr} + v'_{dr} \, i'_{dr} + 2v'_{0r} \, i'_{0r} \right)$$

• Substituting the voltage equations in the above expression yields  $P_{in} = \frac{3}{2} \left( r_s \, i_{qs}^2 + i_{qs} \, p \lambda_{qs} + \omega \lambda_{ds} i_{qs} + r_s \, i_{ds}^2 + i_{ds} \, p \lambda_{ds} - \omega \lambda_{qs} i_{ds} + 2r_s \, i_{0s}^2 + 2i_{0s} \, p \lambda_{0s} \right)$   $+ r_r' \, i_{qr}'^2 + i_{qr}' \, p \lambda_{qr}' + (\omega - \omega_r) \lambda_{dr}' i_{qr}' + r_r' \, i_{dr}'^2 + i_{dr}' \, p \lambda_{dr}' - (\omega - \omega_r) \lambda_{qr}' i_{dr}' + 2r_r' \, i_{0r}'^2 + 2i_{0r}' \, p \lambda_{0r}' \right)$
### Induction Machine Modelling in Arbitrary qd0 RF Torque Equation

$$P_{in} = \frac{3}{2} \left( r_s \, i_{qs}^2 + i_{qs} \, p \lambda_{qs} + \omega \lambda_{ds} i_{qs} + r_s \, i_{ds}^2 + i_{ds} \, p \lambda_{ds} - \omega \lambda_{qs} i_{ds} + 2r_s \, i_{0s}^2 + 2i_{0s} \, p \lambda_{0s} \right)$$
  
+  $r_r' \, i_{qr}'^2 + i_{qr}' \, p \lambda_{qr}' + (\omega - \omega_r) \lambda_{dr}' i_{qr}' + r_r' \, i_{dr}'^2 + i_{dr}' \, p \lambda_{dr}' - (\omega - \omega_r) \lambda_{qr}' i_{dr}' + 2r_r' \, i_{0r}'^2 + 2i_{0r}' \, p \lambda_{0r}' \right)$ 

- The ohmic losses (*ri*<sup>2</sup> terms) and the rate of exchange of magnetic field energy between windings (*ip*λ terms) do not contribute to electromagnetic (EM) torque development.
- Hence the **power terms contributed** to EM torque generation are:

$$P_T = \frac{3}{2} \Big[ \omega \Big( \lambda_{ds} i_{qs} - \lambda_{qs} i_{ds} \Big) + (\omega - \omega_r) \Big( \lambda'_{dr} i'_{qr} - \lambda'_{qr} i'_{dr} \Big) \Big]$$

### Induction Machine Modelling in Arbitrary qd0 RF Torque Equation

$$P_{T} = \frac{3}{2} \Big[ \omega \Big( \lambda_{ds} i_{qs} - \lambda_{qs} i_{ds} \Big) + (\omega - \omega_{r}) \Big( \lambda_{dr}' i_{qr}' - \lambda_{qr}' i_{dr}' \Big) \Big]$$

• The electromagnetic torque is obtained by dividing the above power by the **mechanical speed**:

$$T_{em} = \frac{3}{2} \frac{P}{2\omega_r} \left[ \omega \left( \lambda_{ds} i_{qs} - \lambda_{qs} i_{ds} \right) + (\omega - \omega_r) \left( \lambda'_{dr} i'_{qr} - \lambda'_{qr} i'_{dr} \right) \right]$$

where *P* is the number of poles.

• It can be shown that:

$$\lambda_{ds}i_{qs} - \lambda_{qs}i_{ds} = -(\lambda'_{dr}i'_{qr} - \lambda'_{qr}i'_{dr}) = L_m(i'_{dr}i_{qs} - i'_{qr}i_{ds})$$

### Induction Machine Modelling in Arbitrary qd0 RF Torque Equation

$$T_{em} = \frac{3}{2} \frac{P}{2\omega_r} \left[ \omega \left( \lambda_{ds} i_{qs} - \lambda_{qs} i_{ds} \right) + (\omega - \omega_r) \left( \lambda'_{dr} i'_{qr} - \lambda'_{qr} i'_{dr} \right) \right]$$

$$\lambda_{ds}i_{qs} - \lambda_{qs}i_{ds} = -(\lambda'_{dr}i'_{qr} - \lambda'_{qr}i'_{dr}) = L_m(i'_{dr}i_{qs} - i'_{qr}i_{ds})$$

• Using the above relations, the following expressions for the electromagnetic torque calculation are obtained:

$$T_{em} = \frac{3}{2} \frac{P}{2} \left( \lambda'_{qr} i'_{dr} - \lambda'_{dr} i'_{qr} \right)$$
$$T_{em} = \frac{3}{2} \frac{P}{2} \left( \lambda_{ds} i_{qs} - \lambda_{qs} i_{ds} \right)$$
$$T_{em} = \frac{3}{2} \frac{P}{2} L_m \left( i'_{dr} i_{qs} - i'_{qr} i_{ds} \right)$$

### Induction Machine Modelling in Arbitrary qd0 RF Mechanical Dynamic Equation

• Based on Newton's 2<sup>nd</sup> law for rotational movement,  $\sum T = J \frac{d\omega_{rm}}{dt}$  the **mechanical dynamic equation** is obtained:

**Motoring Mode** 

**Generating Mode** 

$$T_{em} - T_{mech} - T_{damp} = J \frac{d\omega_{rm}}{dt}$$

$$T_{em} + T_{mech} - T_{damp} = J \frac{d\omega_{rm}}{dt}$$

where  $T_{mech}$  is the externally-applied mechanical torque,  $T_{damp}$  is the damping torque, J is the moment of inertia and  $\omega_{rm}$  is the mechanical rotational velocity.

## Induction Machine Modelling in qd0 RF



 From the dynamic equations in *qd*0 reference frame with arbitrary rotating speed, the equations in the same reference frame with other rotating speed can be obtained:





#### Arbitrary rotating speed $\omega$





#### Stationary Reference Frame $\omega = 0$





**Fixed-to-rotor Reference Frame**  $\omega = \omega_r$ 





**Synchronous Reference Frame**  $\omega = \omega_e$ 



$$T_{em} = \frac{3}{2} \frac{P}{2\omega_b} \left( \psi'_{qr} i'_{dr} - \psi'_{dr} i'_{qr} \right) = \frac{3}{2} \frac{P}{2\omega_b} \left( \psi_{ds} i_{qs} - \psi_{qs} i_{ds} \right) = \frac{3}{2} \frac{P}{2\omega_b} x_m \left( i'_{dr} i_{qs} - i'_{qr} i_{ds} \right)$$

• To simulate the induction machines the inputs, outputs and states should be defined





### Sub-systems: Transformation from *abc* to *qd*0





#### Sub-systems: Voltage & flux linkage equations



The stator and rotor equations are coupled.



**Sub-systems: Torque equation** 



Note that rotor quantities or a combination of rotor and stator quantities can be used for electromagnetic torque calculation.



### **Sub-systems: Mechanical equation**





### **Sub-systems**: Transformation from *qd*0 to *abc*



**Voltage sources to windings connections** 





- The 3 applied voltages to the stator terminals need not be balanced nor sinusoidal.
- The 3 stator phase voltages are



3-phase P-pole symmetrical IM

•  $v_{sg}$  can be determined as

$$v_{sg} = R_{sg}(i_{as} + i_{bs} + i_{cs}) + L_{sg} \frac{d}{dt}(i_{as} + i_{bs} + i_{cs}) = 3\left(R_{sg} + L_{sg} \frac{d}{dt}\right) i_{0s}$$

where  $R_{sg}$  and  $L_{sg}$  are the resistance and inductance of the connection between the two neutral points, s and g.

- Clearly when s and g are connected directly v<sub>sg</sub> = 0.
- In the above expression if  $i_{0s}$ = 0 then  $v_{sg}$  = 0.





3-phase P-pole symmetrical IM

- In the case of 4-wire connection (neutral points are connected together), i<sub>0s</sub> is zero if a symmetrical induction machine is supplied by a set of applied voltages that are sinusoidal and balanced.
- In the case of 3-wire connection, i<sub>0s</sub> is always zero by physical constraint, irrespective of whether the 3-phase currents are balanced or not.
- In the case of 3-wire connection, although i<sub>0s</sub> is always zero, v<sub>sg</sub> may not be zero depending on whether the applied voltages are sinusoidal and balanced.



- When the applied voltages are non-sinusoidal (such as the voltages of a six-step inverter), the zero-sequence component of the applied voltage may not be zero.
- In the case of **3-wire connection**,  $v_{sg}$  can be determined, in simulation, using the following expression with  $L_{sg}$  =0 and  $R_{sg}$  set to a high value to approximate the open-circuit condition.

$$v_{sg} = R_{sg}(i_{as} + i_{bs} + i_{cs}) + L_{sg} \frac{d}{dt}(i_{as} + i_{bs} + i_{cs}) = 3\left(R_{sg} + L_{sg} \frac{d}{dt}\right)i_{0s}$$
  
High value 0



Transformation of the stator voltages from stator *abc* system to the stationary qd0 RF where q-axis is aligned with the stator a-axis.

$$\begin{bmatrix} \mathbf{T}_{qd0} \end{bmatrix} = \frac{2}{3} \begin{bmatrix} \cos\theta & \cos\left(\theta - \frac{2\pi}{3}\right) & \cos\left(\theta + \frac{2\pi}{3}\right) \\ \sin\theta & \sin\left(\theta - \frac{2\pi}{3}\right) & \sin\left(\theta + \frac{2\pi}{3}\right) \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix} \xrightarrow{\theta = 0} \begin{bmatrix} \mathbf{T}_{qd0} \end{bmatrix} = \frac{2}{3} \begin{bmatrix} 1 & -\frac{1}{2} & -\frac{1}{2} \\ 0 & -\frac{\sqrt{3}}{2} & \frac{\sqrt{3}}{2} \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix}$$

$$\begin{bmatrix} \mathbf{v}_{qd0} \end{bmatrix} = \begin{bmatrix} \mathbf{T}_{qd0} \end{bmatrix} \begin{bmatrix} \mathbf{v}_{abc} \end{bmatrix}$$

$$\begin{bmatrix} v_{qs}^{s} = \frac{2}{3} v_{as} - \frac{1}{3} v_{bs} - \frac{1}{3} v_{cs} = \frac{2}{3} v_{ag} - \frac{1}{3} v_{bg} - \frac{1}{3} v_{cg} \\ v_{ds}^{s} = \frac{-1}{\sqrt{3}} (v_{bs} - v_{cs}) = \frac{-1}{\sqrt{3}} (v_{bg} - v_{cg}) \\ v_{0s}^{s} = \frac{1}{3} (v_{as} + v_{bs} + v_{cs}) = \frac{1}{3} (v_{ag} + v_{bg} + v_{cg}) - v_{sg} \end{bmatrix}$$

$$93$$



Transformation of the rotor voltages referred to stator side from rotor *abc* system to the *qd*0 RF fixed to rotor where *q*-axis is aligned with the rotor *a*-axis.

$$v_{qr}^{\prime r} = \frac{2}{3}v_{ar}^{\prime} - \frac{1}{3}v_{br}^{\prime} - \frac{1}{3}v_{cr}^{\prime} = \frac{2}{3}v_{an}^{\prime} - \frac{1}{3}v_{bn}^{\prime} - \frac{1}{3}v_{cn}^{\prime}$$

$$v_{dr}^{\prime r} = \frac{-1}{\sqrt{3}}(v_{br}^{\prime} - v_{cr}^{\prime}) = \frac{-1}{\sqrt{3}}(v_{bn}^{\prime} - v_{cn}^{\prime})$$

$$v_{0r}^{\prime r} = \frac{1}{3}(v_{ar}^{\prime} + v_{br}^{\prime} + v_{cr}^{\prime}) = \frac{1}{3}(v_{an}^{\prime} + v_{bn}^{\prime} + v_{cn}^{\prime}) - v_{rn}^{\prime}$$
Rotor winding

where  $v'_{rn}$  is the voltage between points r and n. The prime denotes values referred to the stator side.

**Transformation (rotor voltage; 2<sup>nd</sup> transformation)** 

• Transformation of the **rotor voltages referred to stator side** from the *qd*0 RF **fixed to rotor** to the **stationary** *qd*0 RF.

$$v_{qr}^{\prime s} = v_{qr}^{\prime r} \cos \theta_r + v_{dr}^{\prime r} \sin \theta_r$$

$$v_{dr}^{\prime s} = -v_{qr}^{\prime r} \sin \theta_r + v_{dr}^{\prime r} \cos \theta_r$$
where the rotor angle is
$$\theta_r = \int_0^t \omega_r \, dt + \theta_r(0)$$

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×

q-axis flux linkage equations

stator 
$$\psi_{qs}^{s} = x_{ls} i_{qs}^{s} + \psi_{mq}^{s} \longrightarrow i_{qs}^{s} = \frac{\psi_{qs}^{s} - \psi_{mq}^{s}}{x_{ls}} \mathbf{1}$$
  
Nagnetiting  $\psi_{mq}^{s} = x_{m} (i_{qs}^{s} + i_{qr}^{s})$   
 $\psi_{mq}^{s} = x_{m} (i_{qs}^{s} + i_{qr}^{s})$   
 $\psi_{mq}^{s} = x_{m} (\frac{\psi_{qs}^{s}}{x_{ls}} + \frac{\psi_{qr}^{s}}{x_{lr}})$   
where  
 $\frac{1}{x_{M}} = \frac{1}{x_{m}} + \frac{1}{x_{ls}} + \frac{1}{x_{lr}'}$   
 $g_{gs}^{s} = \frac{\psi_{qr}^{s} - \psi_{mq}^{s}}{x_{lr}'}$ 

d-axis flux linkage equations

stator 
$$\psi_{ds}^{s} = x_{ls} i_{ds}^{s} + \psi_{md}^{s}$$
  $\Longrightarrow$   $i_{ds}^{s} = \frac{\psi_{ds}^{s} - \psi_{md}^{s}}{x_{ls}}$  6  
 $\psi_{ds}^{s} = x_{ls} i_{ds}^{s} + \psi_{md}^{s}$   $\psi_{md}^{s} = x_{m} (i_{ds}^{s} + i_{dr}^{\prime s})$   $\psi_{md}^{s} = x_{M} \left( \frac{\psi_{ds}^{s}}{x_{ls}} + \frac{\psi_{dr}^{\prime s}}{x_{lr}^{\prime}} \right)$  8  
where  $\psi_{dr}^{\prime s} = x_{lr}^{\prime} i_{dr}^{\prime s} + \psi_{md}^{s}$   $\longleftrightarrow$   $i_{dr}^{\prime s} = \frac{\psi_{dr}^{\prime s} - \psi_{md}^{s}}{x_{lr}^{\prime}}$  7

#### **Simulation of Induction Machines** in the Stationary RF d-axis **Stator voltage equations** $\frac{d\psi_{qs}^{s}}{dt}$ $i_{qs}^s = \frac{\psi_{qs}^s - \psi_{qs}^s}{1-1}$ $\psi_{\underline{mq}}$ & $\psi_{qs}^{s} = \omega_{b} \int \left\{ v_{qs}^{s} + \frac{r_{s}}{x_{bs}} \left( \psi_{mq}^{s} - \psi_{qs}^{s} \right) \right\}$ $v_{qs}^s = r_s \, i_{qs}^s + \frac{1}{\omega_b}$ dt dt d-anis $\frac{d\psi_{ds}^{s}}{dt} \& i_{ds}^{s} = \frac{\psi_{ds}^{s} - \psi_{md}^{s}}{x_{ls}}$ $v_{ds}^s = r_s \, i_{ds}^s + \frac{1}{\omega_b} \frac{c}{\omega_b}$ $\psi_{ds}^{s} = \omega_{b} \int \left\{ v_{ds}^{s} + \frac{r_{s}}{r_{b}} \left( \psi_{md}^{s} - \psi_{ds}^{s} \right) \right\}$ 1ero $v_{0s} = r_s i_{0s} + \frac{1}{\omega_b} \frac{d\psi_{0s}}{dt}$ $i_{0s} = \frac{\omega_b}{x_{ls}} \int (v_{0s} - r_s i_{0s}) dt$ **&** $\psi_{0s} = x_{ls} i_{0s}$



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q-axis voltage and flux linkage block diagram



d-axis voltage and flux linkage block diagram



### **Torque and Mechanical equations**

• The torque equation is

$$T_{em} = \frac{3}{2} \frac{P}{2\omega_b} \left( \psi_{ds} i_{qs} - \psi_{qs} i_{ds} \right)$$

• The dynamic mechanical equation is

Motoring Mode
$$T_{em} - T_{mech} - T_{damp} = J \frac{d\omega_{rm}}{dt}$$
Generating Mode $T_{em} + T_{mech} - T_{damp} = J \frac{d\omega_{rm}}{dt}$ 



• Transformation of the stator currents from the stationary qd0 RF to the stator *abc* system.

 $\begin{bmatrix} \mathbf{T}_{qd0} \end{bmatrix}^{-1} = \begin{vmatrix} -\frac{1}{2} & -\frac{\sqrt{3}}{2} & 1 \\ -\frac{1}{2} & \frac{\sqrt{3}}{2} & 1 \end{vmatrix}$ 

$$[\mathbf{i}_{abc}] = [\mathbf{T}_{qd0}]^{-1} [\mathbf{i}_{qd0}] \text{ where}$$
$$i_{as} = i_{qs}^{s} + i_{0s}$$
$$i_{bs} = -\frac{1}{2} i_{qs}^{s} - \frac{\sqrt{3}}{2} i_{ds}^{s} + i_{0s}$$
$$i_{cs} = -\frac{1}{2} i_{qs}^{s} + \frac{\sqrt{3}}{2} i_{ds}^{s} + i_{0s}$$



• Transformation of the **rotor currents referred to stator side** from the **stationary** *qd*0 RF to the *qd*0 RF **fixed to rotor**.

$$i_{qr}^{\prime r} = i_{qr}^{\prime s} \cos \theta_r - i_{dr}^{\prime s} \sin \theta_r$$
$$i_{dr}^{\prime r} = i_{qr}^{\prime s} \sin \theta_r + i_{dr}^{\prime s} \cos \theta_r$$

where the rotor angle is

$$\theta_r = \int_0^t \omega_r \, dt + \theta_r(0)$$



**Transformation (rotor currents; 2<sup>nd</sup> transformation)** 

• Transformation of the **rotor currents referred to stator side** from the *qd*0 RF **fixed to rotor** to the rotor *abc* system.



## Per-Unit System



#### **Base quantities**

 The base quantities with peak rather than rms value of a P-pole, three-phase induction machine with rated line-to-line rms voltage, V<sub>rated</sub>, and rated volt-ampere, S<sub>rated</sub>, are as follows:



### **Per-Unit System**



### **Torque Equation**

$$T_{em} = \frac{3}{2} \frac{P}{2\omega_b} \left( \psi_{ds} i_{qs} - \psi_{qs} i_{ds} \right)$$

$$T_{b} = \frac{S_{b}}{\omega_{bm}} \xrightarrow{\omega_{bm}} T_{b} = \frac{S_{b}}{\frac{2}{p}\omega_{b}} \xrightarrow{I_{b}} T_{b} = \frac{S_{b}}{\frac{2}{p}\omega_{b}} \xrightarrow{I_{b}} T_{b} = \frac{S_{b}}{\frac{2}{p}\omega_{b}} \xrightarrow{I_{b}} T_{b} = \frac{3}{2} \frac{P}{2} \frac{V_{b}I_{b}}{\omega_{b}}$$

$$T_{b} = \frac{\frac{3}{2} \frac{P}{2\omega_{b}} (\psi_{ds}i_{qs} - \psi_{qs}i_{ds})}{\frac{3}{2} \frac{P}{2\omega_{b}} V_{b}I_{b}} \xrightarrow{T_{em}(pu)} = \psi_{ds}(pu)i_{qs}(pu) - \psi_{qs}(pu)i_{ds}(pu)$$

#### where

$$\psi_{ds}(pu) = \frac{\psi_{ds}}{V_b} \qquad \psi_{qs}(pu) = \frac{\psi_{qs}}{V_b} \qquad i_{qs}(pu) = \frac{i_{qs}}{I_b} \qquad i_{ds}(pu) = \frac{i_{ds}}{I_b} \qquad 107$$

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### **Per-Unit System**



### **Mechanical Equation**



where  $H = J\omega_{bm}^2 / (2S_b)$  is the inertia constant.
# **Machine Parameters**



The following parameters/quantities are required for the simulation

- $X_{ls}$  stator or armature winding leakage reactance
- $x'_{lr}$  rotor winding leakage reactance referred to stator
- $x_m$  magnetizing reactance
- $r_s$  stator or armature winding resistance
- $r'_r$  rotor winding resistance referred to stator
- J rotor moment of inertia
- *P* number of poles
- $\omega_b$  rotor base speed
- $v_{rated}$  rated line-to-line rms voltage of stator
- *S<sub>rared</sub>* rated **volt-ampere**
- $T_{mech}$  mechanical load torque profile



- The iron saturation mainly affects the value of the magnetizing inductance and, to a much lesser extent, the leakage inductances.
- Considering the iron saturation effects on the leakage inductances is complex (because of the complicated leakage flux path) and these effects are neglected here.
- Therefore the iron saturation effects on the **magnetizing inductance** are only **considered**.
- We assume that the iron saturation affects the *q* and *d*-axis components in the **same manner** (due to smooth air-gap).





#### **Saturation Characteristics**



• The saturated value of the mutual flux linkage per second in the *q*-axis is given by:

$$\psi_{mq}^{s,sat} = \psi_{mq}^{s,unsat} - \Delta \psi_{mq}^{s}$$

$$\psi_{mq}^{s,sat} = x_m \left( \frac{\psi_{qs}^s - \psi_{mq}^{s,sat}}{x_{ls}} + \frac{\psi_{qr}^{\prime s} - \psi_{mq}^{s,sat}}{x_{lr}^{\prime}} \right) - \Delta \psi_{mq}^s$$

• Rearranging subject to  $\Psi_{mq}^{s,sat}$  yields

$$\psi_{mq}^{s,sat} = x_M \left( \frac{\psi_{qs}^s}{x_{ls}} + \frac{\psi_{qr}^{\prime s}}{x_{lr}^{\prime}} - \frac{\Delta \psi_{mq}^s}{x_m} \right) \quad \text{where} \quad \left[ \frac{1}{x_M} \right]$$

Unsaturated  
$$\frac{1}{x_M} = \frac{1}{x_m} + \frac{1}{x_{ls}} + \frac{1}{x'_{lr}}$$



• Similarly the saturated value of the mutual flux linkage per second in the *d*-axis is given by:

 $\psi_{md}^{s,sat} = \psi_{md}^{s,unsat} - \Delta \psi_{md}^{s}$ 

$$\psi_{md}^{s,sat} = x_m \left( \frac{\psi_{ds}^s - \psi_{md}^{s,sat}}{x_{ls}} + \frac{\psi_{dr}^{\prime s} - \psi_{md}^{s,sat}}{x_{lr}^{\prime}} \right) - \Delta \psi_{md}^s$$

• Rearranging subject to  $\Psi_{md}^{s,sat}$  yields

$$\psi_{md}^{s,sat} = x_M \left( \frac{\psi_{ds}^s}{x_{ls}} + \frac{\psi_{dr}^{\prime s}}{x_{lr}^{\prime}} - \frac{\Delta \psi_{md}^s}{x_m} \right) \quad \text{where} \quad \left| \frac{1}{x_M} \right| =$$

Unsaturated  
$$\frac{1}{x_M} = \frac{1}{x_m} + \frac{1}{x_{ls}} + \frac{1}{x'_{lr}}$$



• Assuming a **proportional reduction** in flux linkages of the *q*- and *d*-axis yields:



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- The relationship between  $\Delta \psi_m$  and  $\psi_m^{sat}$  can be determined from the no-load test curve of the machine.
- The value of  $\Delta \psi_m$  is obtained from  $\psi_m^{sat}$  using piece-wise segments or a look-up table as explained for the transformers.







• The part of the IM simulation that is **affected** by the **inclusion of mutual flux saturation** is shown below:





• For nonlinear models, the well-established linear control techniques cannot be employed.

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{f}(\mathbf{x}, \mathbf{u}) \\ \mathbf{y} = \mathbf{g}(\mathbf{x}, \mathbf{u}) \end{cases} \longrightarrow \text{Linear control theory cannot be used.}$$

• If the nonlinear model can be linearized around an equilibrium point, then the linear control techniques **can** be used.



#### When can a nonlinear model be linearized?

- If the system is always working around an equilibrium point.
   Or
- 2. If the nonlinear terms are **not significant** and can be approximated by linear terms.

#### The linearization methods

- **1.** Taylor series expansion
- 2. Perturbation method

Nonlinear IM model in the arbitrary qd0 RF



#### Nonlinear IM model in the synchronous qd0 RF



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# **First Technique:**

# **Taylor Series Expansion**

Consider a system whose input is x(t) and output is y(t). The relationship between y(t) and x(t) is given by v = f(x)

If the normal operating condition corresponds to  $x_o$ ,  $y_o$  then above relation may be expanded into a Taylor series about this point as follows:  $\frac{df}{dt} = \frac{1}{2} \frac{d^2 f}{dt}$ 

$$y = f(x_o) + \frac{df}{dx}\Big|_{x=x_o} (x - x_o) + \frac{1}{2!} \frac{d^2 f}{dx^2}\Big|_{x=x_o} (x - x_o)^2 + \cdots$$

Neglecting the higher-order terms yields  $y = y_o + K(x - x_o)$ 

where 
$$y_o = f(x_o)$$
 and  $K = \frac{df}{dx}\Big|_{x=x_o}$ 

Finally assuming  $\Delta x = x - x_o$  and  $\Delta y = y - y_o$  yields  $\Delta y = K \Delta x$ 







Consider a system whose input is  $x_1(t)$  and  $x_2(t)$  and output is y(t). The relationship between y(t) and the inputs is given by

$$y = f(x_1, x_2)$$

Expanding Taylor series about equilibrium point  $x_{1o}$ ,  $x_{2o}$ ,  $y_o$ 

$$y = f(x_{1o}, x_{2o}) + \left[ \frac{\partial f}{\partial x_1} \Big|_{\substack{x_1 = x_{1o} \\ x_2 = x_{2o}}} (x_1 - x_{1o}) + \frac{\partial f}{\partial x_2} \Big|_{\substack{x_1 = x_{1o} \\ x_2 = x_{2o}}} (x_2 - x_{2o}) \right]$$
$$+ \frac{1}{2!} \left[ \frac{\partial^2 f}{\partial x_1^2} \Big|_{\substack{x_1 = x_{1o} \\ x_2 = x_{2o}}} (x_1 - x_{1o})^2 + \frac{\partial^2 f}{\partial x_2^2} \Big|_{\substack{x_1 = x_{1o} \\ x_2 = x_{2o}}} (x_2 - x_{2o})^2 \right]$$
$$+ 2 \frac{\partial^2 f}{\partial x_1 \partial x_2} \Big|_{\substack{x_1 = x_{1o} \\ x_2 = x_{2o}}} (x_1 - x_{1o}) (x_2 - x_{2o}) \Big] + \cdots$$



Neglecting the higher-order terms yields

$$y = f(x_1, x_2) \implies y = f(x_{1o}, x_{2o}) + \left[\frac{\partial f}{\partial x_1}\Big|_{\substack{x_1 = x_{1o} \\ x_2 = x_{2o}}} (x_1 - x_{1o}) + \frac{\partial f}{\partial x_2}\Big|_{\substack{x_1 = x_{1o} \\ x_2 = x_{2o}}} (x_2 - x_{2o})\right]$$

which can be written as 
$$y = y_o + K_1(x_1 - x_{1o}) + K_2(x_2 - x_{2o})$$
  
where  $y_o = f(x_{1o}, x_{2o})$ ,  $K_1 = \frac{\partial f}{\partial x_1}\Big|_{\substack{x_1 = x_{1o} \\ x_2 = x_{2o}}}$  and  $K_2 = \frac{\partial f}{\partial x_2}\Big|_{\substack{x_1 = x_{1o} \\ x_2 = x_{2o}}}$ 

Finally assuming  $\Delta x_1 = x_1 - x_{1o}$ ,  $\Delta x_2 = x_2 - x_{2o}$  and  $\Delta y = y - y_o$  yields

$$\Delta y = K_1 \,\Delta x_1 + K_2 \,\Delta x_2$$



Linearization of the system  $y = f(x_1, \dots, x_n)$  yields

$$\Rightarrow \qquad y = f(x_{1o}, \dots, x_{no}) + \left[ \frac{\partial f}{\partial x_1} \Big|_{\substack{x_1 = x_{1o} \\ \vdots \\ x_n = x_{no}}} (x_1 - x_{1o}) + \dots + \frac{\partial f}{\partial x_n} \Big|_{\substack{x_1 = x_{1o} \\ \vdots \\ x_n = x_{no}}} (x_n - x_{no}) \right]$$
  
which can be written as 
$$y = y_o + K_1(x_1 - x_{1o}) + \dots + K_n(x_n - x_{no})$$
  
where  $y_o = f(x_{1o}, \dots, x_{no})$ ,  $K_1 = \frac{\partial f}{\partial x_1} \Big|_{\substack{x_1 = x_{1o} \\ \vdots \\ x_n = x_{no}}} \text{ and } K_n = \frac{\partial f}{\partial x_n} \Big|_{\substack{x_1 = x_{1o} \\ \vdots \\ x_n = x_{no}}} x_n = x_{no}$ 

Finally assuming  $\Delta x_1 = x_1 - x_{1o}$ ,  $\Delta x_n = x_n - x_{no}$  and  $\Delta y = y - y_o$  yields

$$\Delta y = K_1 \,\Delta x_1 + \dots + K_n \,\Delta x_n$$



**Example:** Linearize the following nonlinear algebraic equation about  $x_{1o} = 2$  and  $x_{2o} = 1$ 

$$y = 2x_1^2 x_2 + \sin(\pi x_1) + \sqrt{2x_1 x_2}$$

**Solution**:



Linearization of Nonlinear Systems  
Solution: 
$$x_{1o} = 2$$
  $x_{2o} = 1$   $y = 2x_1^2 x_2 + \sin(\pi x_1) + \sqrt{2x_1 x_2}$   
 $\Delta y = K_1 \Delta x_1 + K_2 \Delta x_2$   $\Delta y = (8.5 + \pi)\Delta x_1 + 9\Delta x_2$   
 $y_o = f(x_{1o}, x_{2o}) = 2 \times 2^2 \times 1 + \sin(2\pi) + \sqrt{2 \times 2 \times 1} = 10$   
 $\Delta y = y - y_o = y - 10$   $\Delta x_1 = x_1 - x_{1o} = x_1 - 2$   $\Delta x_2 = x_2 - x_{2o} = x_2 - 1$   
 $K_1 = \frac{\partial f}{\partial x_1}\Big|_{\substack{x_1 = x_{1o} \\ x_2 = x_{2o}}} = \left(4x_1 x_2 + \pi \cos(\pi x_1) + \frac{x_2}{\sqrt{2x_1 x_2}}\right)\Big|_{\substack{x_1 = 2 \\ x_2 = 1}} = 8 + \pi + \frac{1}{2} = (8.5 + \pi)$   
 $K_2 = \frac{\partial f}{\partial x_2}\Big|_{\substack{x_1 = x_{1o} \\ x_2 = x_{2o}}} = \left(2x_1^2 + \frac{x_1}{\sqrt{2x_1 x_2}}\right)\Big|_{\substack{x_1 = 2 \\ x_2 = 1}} = 8 + 1 = 9$   
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• Consider the following operating point

 $P_o = \begin{pmatrix} i^e_{qs,o} & i^e_{ds,o} & i'^e_{qr,o} & i'^e_{dr,o} \end{pmatrix}$ 

• It is required to linearize the following EM torque expression around the given operating point

$$T_{em} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left( i'_{dr} i_{qs} - i'_{qr} i_{ds} \right)$$

• The EM torque at the operating point is

$$T_{em} \mid_{P_o} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left( i'_{dr,o} \ i_{qs,o} - i'_{qr,o} \ i_{ds,o} \right)$$

Linearization of electromagnetic torque expression

$$T_{em} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left( i'_{dr} i_{qs} - i'_{qr} i_{ds} \right)$$

$$P_o = ig( egin{array}{cccc} i^e_{qs,o} & i^e_{ds,o} & i^{\prime e}_{qr,o} & i^{\prime e}_{dr,o} ig) \end{array}$$

$$T_{em}|_{P_o} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left( i'_{dr,o} i_{qs,o} - i'_{qr,o} i_{ds,o} \right)$$

• Using the Taylor series expansion yields

$$T_{em} = T_{em} \mid_{P_o} + \left[ \frac{\partial T_{em}}{\partial i_{qs}^e} \right|_{P_o} \Delta i_{qs}^e + \frac{\partial T_{em}}{\partial i_{ds}^e} \right|_{P_o} \Delta i_{ds}^e + \frac{\partial T_{em}}{\partial i_{ds}^{\prime e}} \right|_{P_o} \Delta i_{qr}^{\prime e} + \frac{\partial T_{em}}{\partial i_{qr}^{\prime e}} \right|_{P_o} \Delta i_{dr}^{\prime e} + \frac{\partial T_{em}}{\partial i_{dr}^{\prime e}} \right|_{P_o} \Delta i_{dr}^{\prime e} = 0$$

$$\Delta T_{em} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left[ i_{dr,o}^{\prime e} \Delta i_{qs}^{e} - i_{qr,o}^{\prime e} \Delta i_{ds}^{e} - i_{ds,o}^{e} \Delta i_{qr}^{\prime e} + i_{qs,o}^{e} \Delta i_{dr}^{\prime e} \right]$$

#### Linearization of electromagnetic torque expression

$$\Delta T_{em} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left[ i_{dr,o}^{\prime e} \Delta i_{qs}^{e} - i_{qr,o}^{\prime e} \Delta i_{ds}^{e} - i_{ds,o}^{e} \Delta i_{qr}^{\prime e} + i_{qs,o}^{e} \Delta i_{dr}^{\prime e} \right]$$

where

$$\Delta T_{em} = T_{em} - T_{em} \mid_{P_o} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left( i'_{dr,o} i_{qs,o} - i'_{qr,o} i_{ds,o} \right)$$

$$\Delta i_{qs}^e = i_{qs}^e - i_{qs,o}^e \qquad \Delta i_{ds}^e = i_{ds}^e - i_{ds,o}^e$$

$$\Delta i_{qr}^{\prime e} = i_{qr}^{\prime e} - i_{qr,o}^{\prime e} \qquad \Delta i_{dr}^{\prime e} = i_{dr}^{\prime e} - i_{dr,o}^{\prime e}$$





# **Second Technique:**

# **Perturbation Method**



• Consider a **nonlinear model** in the following form

$$\begin{cases} \mathbf{f}(\dot{\mathbf{x}}, \mathbf{x}, \mathbf{u}) = 0\\ \mathbf{y} = \mathbf{g}(\mathbf{x}, \mathbf{u}) \end{cases}$$

where  $\mathbf{x}$  is the vector of state variables,  $\mathbf{u}$  is the input vector and  $\mathbf{y}$  is the output vector.

- Assume the system is operating in an equilibrium point in which x=x<sub>o</sub>, u=u<sub>o</sub> and y=y<sub>o</sub>.
- When a small displacement, denoted by ∆, is applied to each component of x, u and y vectors, the perturbed variables will still satisfy the governing differential-algebraic equations:

$$\begin{cases} \mathbf{f} \left( \dot{\mathbf{x}}_{\mathbf{x}=\mathbf{x}_{o}} + \Delta \dot{\mathbf{x}}, \ \mathbf{x}_{o} + \Delta \mathbf{x}, \ \mathbf{u}_{o} + \Delta \mathbf{u} \right) = 0 \\ \mathbf{y}_{o} + \Delta \mathbf{y} = \mathbf{g} \left( \mathbf{x}_{o} + \Delta \mathbf{x}, \ \mathbf{u}_{o} + \Delta \mathbf{u} \right) \end{cases}$$



$$\begin{cases} \mathbf{f} \left( \dot{\mathbf{x}}_{\mathbf{x}=\mathbf{x}_{o}} + \Delta \dot{\mathbf{x}}, \mathbf{x}_{o} + \Delta \mathbf{x}, \mathbf{u}_{o} + \Delta \mathbf{u} \right) = 0 \\ \mathbf{y}_{o} + \Delta \mathbf{y} = \mathbf{g} \left( \mathbf{x}_{o} + \Delta \mathbf{x}, \mathbf{u}_{o} + \Delta \mathbf{u} \right) \end{cases} = 0$$

• At the equilibrium point we have

 Substituting these two relations into the perturbed state-space equations and neglecting higher order ∆ terms can result in the following linear state-space equations:

$$\begin{cases} \Delta \dot{\mathbf{x}} = \mathbf{A} \Delta \mathbf{x} + \mathbf{B} \Delta \mathbf{u} \\ \Delta \mathbf{y} = \mathbf{C} \Delta \mathbf{x} + \mathbf{D} \Delta \mathbf{u} \end{cases}$$



• Consider the following operating point

 $P_o = \begin{pmatrix} i_{qs,o}^e & i_{ds,o}^e & i_{qr,o}^{\prime e} & i_{dr,o}^{\prime e} \end{pmatrix}$ 

• It is required to linearize the following EM torque expression around the given operating point

 $T_{em} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left( i'_{dr} i_{qs} - i'_{qr} i_{ds} \right)$ 

• The EM torque at the operating point is

$$T_{em}|_{P_o} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left( i'_{dr,o} i_{qs,o} - i'_{qr,o} i_{ds,o} \right)$$



Linearization of electromagnetic torque expression

$$T_{em} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left( i'_{dr} i_{qs} - i'_{qr} i_{ds} \right)$$

$$P_o = \begin{pmatrix} i_{qs,o}^e & i_{ds,o}^e & i_{qr,o}^{\prime e} & i_{dr,o}^{\prime e} \end{pmatrix}$$

$$T_{em}|_{P_o} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left( i'_{dr,o} i_{qs,o} - i'_{qr,o} i_{ds,o} \right)$$

• Using the perturbation method yields

$$T_{em}|_{P_o} + \Delta T_{em} = \frac{3}{2} \frac{P}{2\omega_b} x_m \Big[ (i'_{dr,o} + \Delta i'_{dr}) (i_{qs,o} + \Delta i_{qs}) - (i'_{qr,o} + \Delta i'_{qr}) (i_{ds,o} + \Delta i_{ds}) \Big]$$

$$\Delta T_{em} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left[ i_{dr,o}^{\prime e} \Delta i_{qs}^{e} - i_{qr,o}^{\prime e} \Delta i_{ds}^{e} - i_{ds,o}^{e} \Delta i_{qr}^{\prime e} + i_{qs,o}^{e} \Delta i_{dr}^{\prime e} \right]$$

• By using the **perturbation** method, **linearize** the following nonlinear IM model which is in the synchronous *qd*0 RF



Nonlinear terms are encircled.

$$T_{em} - T_{mech} = \frac{2J}{P} \frac{d\omega_r}{dt} \qquad T_{em} = \frac{3}{2} \frac{P}{2\omega_b} x_m \left( (i'_{dr} i_{qs}) - (i'_{qr} i_{ds}) \right) \qquad s = \frac{\omega_e - \omega_r}{\omega_e}$$

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$$\mathbf{x} = \begin{bmatrix} i_{qs}^{e} & i_{ds}^{e} & i_{qr}^{\prime e} & i_{dr}^{\prime e} & \boldsymbol{\omega}_{r} \end{bmatrix}^{T} \qquad \mathbf{u} = \begin{bmatrix} v_{qs}^{e} & v_{ds}^{e} & v_{qr}^{\prime e} & v_{dr}^{\prime e} & T_{mech} \end{bmatrix}^{T}$$

• The linearization is performed around an equilibrium point:

$$\mathbf{x}_{o} = \begin{bmatrix} i_{qs,o}^{e} & i_{ds,o}^{e} & i_{qr,o}^{\prime e} & i_{dr,o}^{\prime e} & \boldsymbol{\omega}_{r,o} \end{bmatrix}^{T} \qquad \mathbf{u}_{o} = \begin{bmatrix} v_{qs,o}^{e} & v_{ds,o}^{e} & v_{qr,o}^{\prime e} & v_{dr,o}^{\prime e} & T_{mech,o} \end{bmatrix}^{T}$$

• The small displacement is introduced for each variables:

$$\Delta \mathbf{x} = \begin{bmatrix} \Delta i_{qs}^{e} & \Delta i_{ds}^{e} & \Delta i_{qr}^{\prime e} & \Delta i_{dr}^{\prime e} & \Delta \omega_{r} \end{bmatrix}^{T} \qquad \Delta \mathbf{u} = \begin{bmatrix} \Delta v_{qs}^{e} & \Delta v_{ds}^{e} & \Delta v_{qr}^{\prime e} & \Delta v_{dr}^{\prime e} & \Delta T_{mech} \end{bmatrix}^{T}$$
where e.g. 
$$\Delta i_{qs}^{e} = i_{qs}^{e} - i_{qs,o}^{e} \qquad \Delta i_{ds}^{e} = i_{ds}^{e} - i_{ds,o}^{e}$$
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- Perturbing the governing equations and neglecting the higher order  $\Delta$  terms yields

$$\begin{bmatrix} \Delta v_{qs}^{e} \\ \Delta v_{ds}^{e} \\ \Delta v_{ds}^{e} \\ \Delta v_{qr}^{e} \\ \Delta v_{dr}^{r} \\ \Delta v_{dr}^{r} \\ \Delta r_{mech}^{r} \end{bmatrix} = \begin{bmatrix} r_{s} + \frac{p}{\omega_{b}} x_{ss} & \frac{\omega_{e}}{\omega_{b}} x_{ss} & \frac{p}{\omega_{b}} x_{m} & \frac{\omega_{e}}{\omega_{b}} x_{m} & 0 \\ -\frac{\omega_{e}}{\omega_{b}} x_{ss} & r_{s} + \frac{p}{\omega_{b}} x_{ss} & -\frac{\omega_{e}}{\omega_{b}} x_{m} & \frac{p}{\omega_{b}} x_{m} & 0 \\ \frac{p}{\omega_{b}} x_{m} & s_{o} \frac{\omega_{e}}{\omega_{b}} x_{m} & r'_{r} + \frac{p}{\omega_{b}} x'_{rr} & s_{o} \frac{\omega_{e}}{\omega_{b}} x'_{rr} & -\frac{x_{m}}{\omega_{b}} i_{ds,o} -\frac{x'_{rr}}{\omega_{b}} i'_{dr,o} \\ -s_{o} \frac{\omega_{e}}{\omega_{b}} x_{m} & \frac{p}{\omega_{b}} x_{m} & -s_{o} \frac{\omega_{e}}{\omega_{b}} x'_{rr} & r'_{r} + \frac{p}{\omega_{b}} x'_{rr} & \frac{x_{m}}{\omega_{b}} i_{qs,o} + \frac{x'_{rr}}{\omega_{b}} i'_{qr,o} \\ \frac{3P}{4\omega_{b}} x_{m} i'_{dr,o} & -\frac{3P}{4\omega_{b}} x_{m} i'_{ds,o} & \frac{3P}{4\omega_{b}} x_{m} i_{qs,o} & -\frac{2J}{P} p \end{bmatrix} \begin{bmatrix} \Delta i_{qs}^{e} \\ \Delta i'_{ds}^{e} \\ \Delta \omega_{r} \end{bmatrix}$$

$$s_o = \frac{\omega_e - \omega_{r,o}}{\omega_e}$$
  $\Delta s = s - s_o = \frac{-\Delta \omega_r}{\omega_e}$ 



 $\Delta \mathbf{u} = \mathbf{E} \Delta \dot{\mathbf{x}} - \mathbf{F} \Delta \mathbf{x}$ 

where

$$\Delta \mathbf{x} = \begin{bmatrix} \Delta i_{qs}^{e} & \Delta i_{ds}^{e} & \Delta i_{qr}^{\prime e} & \Delta i_{dr}^{\prime e} & \Delta \omega_{r} \end{bmatrix}^{T}$$

$$\Delta \mathbf{u} = \begin{bmatrix} \Delta v_{qs}^{e} & \Delta v_{ds}^{e} & \Delta v_{qr}^{\prime e} & \Delta v_{dr}^{\prime e} & \Delta T_{mech} \end{bmatrix}^{T}$$

- and  ${\bf E}$  and  ${\bf F}$  are two constant matrices

$$\mathbf{E} = \frac{1}{\omega_b} \begin{bmatrix} x_{ss} & 0 & x_m & 0 & 0\\ 0 & x_{ss} & 0 & x_m & 0\\ x_m & 0 & x'_{rr} & 0 & 0\\ 0 & x_m & 0 & x'_{rr} & 0\\ 0 & 0 & 0 & 0 & -\frac{2J\omega_b}{P} \end{bmatrix}$$

$$\boldsymbol{F} = -\begin{bmatrix} r_{s} & \frac{\omega_{e}}{\omega_{b}} x_{ss} & 0 & \frac{\omega_{e}}{\omega_{b}} x_{m} & 0 \\ -\frac{\omega_{e}}{\omega_{b}} x_{ss} & r_{s} & -\frac{\omega_{e}}{\omega_{b}} x_{m} & 0 & 0 \\ 0 & s_{o} \frac{\omega_{e}}{\omega_{b}} x_{m} & r_{r}' & s_{o} \frac{\omega_{e}}{\omega_{b}} x_{rr}' & -\frac{x_{m}}{\omega_{b}} i_{ds,o} -\frac{x'_{rr}}{\omega_{b}} i'_{dr,o} \\ -s_{o} \frac{\omega_{e}}{\omega_{b}} x_{m} & 0 & -s_{o} \frac{\omega_{e}}{\omega_{b}} x'_{rr} & r_{r}' & \frac{x_{m}}{\omega_{b}} i_{qs,o} + \frac{x'_{rr}}{\omega_{b}} i'_{qr,o} \\ \frac{3P}{4\omega_{b}} x_{m} i'_{dr,o} & -\frac{3P}{4\omega_{b}} x_{m} i'_{qr,o} & -\frac{3P}{4\omega_{b}} x_{m} i_{ds,o} & \frac{3P}{4\omega_{b}} x_{m} i_{qs,o} & 0 \end{bmatrix}$$

• It is now possible to write the linear state-space equations in the following standard form

 $\Delta \dot{\mathbf{x}} = \mathbf{A} \Delta \mathbf{x} + \mathbf{B} \Delta \mathbf{u}$ 

where 
$$\mathbf{A} = \mathbf{E}^{-1} \mathbf{F}$$
  $\mathbf{B} = \mathbf{E}^{-1}$ 

